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STRUCTURAL IMPLICATIONS OF THE $^{115}\mbox{In}$ NUCLEAR QUADRUPOLE RESONANCE SPECTRUM OF $\mbox{CH}_3\mbox{InI}_2$

93

W.A. WELSH and T.B. BRILL*

Department of Chemistry, University of Delaware, Newark, Delaware 19711 (U.S.A.) (Received October 16th, 1973)

Summary

Based on its vibrational spectrum, CH_3InI_2 has been proposed to contain linear $In(CH_3)_2^*$ ions rather than being a bridged halogen dimer as is found for most other compounds of this type. The ¹¹⁵In nuclear quadrupole resonance spectrum observed here for CH_3InI_2 yields a coupling constant of 1270.6 MHz and an asymmetry parameter of 0.051 ± 0.006 . The small asymmetry parameter is additional evidence of an axially symmetric environment for indium, such as $In(CH_3)_2^*$, and by comparison to In_2I_6 strongly suggests that the dimeric structure is not present.

Recently, the synthesis of CH_3InI_2 by oxidative addition of CH_3I to InI has been reported [1-3]. Unlike other Group IIIA halides and organohalides of this type which contain halogen-bridged dimers or are weakly polymeric in the solid state [4], the ionic formulation, $[In(CH_3)_2]^+[InI_4]^-$, was proposed for CH_3InI_2 . This suggested structure was based on the vibrational spectrum [1-3]. Hobbs and Tobias also reported the existence of $(CH_3)_2In^+$ in aqueous solution [5]. The mass spectrum of CH_3InI_2 studied by Gynane et al. [2], showed the parent peak to be a dimer which would likely have the following structure:

This parent ion was not found by Poland and Tuck [3].

Nuclear quadrupole resonance data could greatly aid in clarifying the structure in the solid state because both ¹¹⁵In (I = 9/2) and ¹²⁷I (I = 5/2) are quadrupolar nuclei and their spectra in the salt form will be very different from those in the bridged dimer form. A complete ¹¹⁵In spectrum would yield four resonances for each type of indium atom present. For ¹²⁷I, two signals will re-

* To whom correspondence should be addressed.

TABLE 1 115 In AND 127 I NUCLEAR QUADRUPOLE RESONANCE DATA (MHz) AND ASYMMETRY PARAM- EMERS IN L. 4 AND CH-ULA AT 25° CP				
115 _{In} 1/2++3/2	3/2++5/2	5/2↔7/2 7/2↔9/2	e ² Qq/h	n
In ₂ I ₆ 37.506 CH ₃ InI ₂ 54.36(10)	29.986 105.55(35)	36.265 51.238 158.68(10) 211.65(5)	320.41 1270.6	0.69 0.051 ± 0.006
127 ₁	1/2↔3/2	3/2⇔5/2	e ² Qq/h	η
In ₂ I ₆ CH ₃ InI ₂	122.728 173.177 173.633	229.190 346.289 347.234	772.43 1154.3 1157.4 —	0.237 ^c 0.0 ^d 0.0 ^d

^a References 6—9. ^b Parenthetical numbers are signal-to-noise ratios. ^c Bridging iodine atom. ^dTerminal iodine atom.

sult from each type. Although very little ¹¹⁵In NQR spectroscopy has been reported, the ¹¹⁵In and ¹²⁷I frequency in In_2I_6 [6–9] and some additional organoindium halide data [10] have been obtained. The nuclear quadrupole coupling constants, $e^2 Qq$, and asymmetry parameters, η , for In_2I_6 are reported in Table 1. With these data and the comparative effects in the analogous AlX₃ and RAIX₂ compounds [11], predictions can be made regarding the spectrum of CH₃InI₂ if the dimeric structure shown above is present in the solid state. Two rather widely separated resonance pairs for iodine corresponding to bridging and terminal iodine atoms should be observed. The lower frequency pair will have a reasonably large asymmetry parameter as a result of the coordinate geometry of the bridging halogen. The terminal halogens will have negligible asymmetry because of their axial symmetry. The indium resonances would show a large asymmetry parameter much like that found in In_2I_6 [6–9] and in the analogous aluminum compounds [11].

If the ionic formulation of $[In(CH_3)_2]^+[InI_4]^-$ is present in the solid state only one set of iodine signals will be observed (assuming no crystallographic inequivalence) and η will be close to zero. For ¹¹⁵In the cation will produce one set of signals yielding ideally zero asymmetry. The anion would give no ¹¹⁵In signals because of the tetrahedral arrangement of iodine atoms about indium [2, 3].

 CH_3InI_2 was prepared as before [1-3]. Colorless highly crystalline material was obtained by dissolving the compound in CH_3I and evaporating the solution over a period of three days using a slow stream of nitrogen. The Raman and mass spectra were recorded and confirmed the compound to be the same as that obtained previously [1-3]. The NQR spectrometer system has been described before [12]. A Wilks OS-4A cavity oscillator was used to scan the 270-350 MHz region.

The results are reported in Table 1. No other signals were found in the 10 - 350 MHz region. The resonances at 105.55, 158.68 and 211.65 MHz are in the region of ¹²⁷I resonances for this type of compound. However the 105.55 and 211.65 MHz signals cannot belong to the same iodine atom because ν (3/2 \leftrightarrow 1/2) must be greater than or equal to one-half ν (5/2 \leftrightarrow 3/2). The difference between

2(105.55) and 211.65 is not great but is at the very outer limit of the experimental uncertainty in the numbers ($\pm 0.08\%$). Further, no high frequency counterpart for the resonance at 158.68 MHz was observed. In fact, it is extremely unlikely that any of these resonances belong to jodine because all four signals fit the eigenvalue equation* for I = 9/2 and $\eta = 0.051 \pm 0.006$ with an uncertainty of 0.08% in the coupling constant. Since the fit of frequencies to the eigenvalue solutions is very sensitive for I = 9/2, all four resonances were determined to belong to ¹¹⁵In. The low asymmetry at indium implies that the bridged formulation represented above and found in RAIX₂ compounds [11] and $\ln_2 I_6$ [6-9] does exist in CH₃InI₂, and that the ionic structure [In(CH₃)₂]⁺ [InI₄]⁻, or a polymeric bridged structure in which the indium atom has axial symmetry as suggested previously [1-3], must be present. The very large coupling constant. (the largest reported to date for ¹¹⁵In), is also consistent with an indium site in which the largest component of the electric field gradient tensors for the external groups adds rather than cancels at In. A linear arrangement for the CH₃ groups about indium would result in such an occurrence.

It would be valuable to observe the ¹²⁷I resonances to have a better understanding of their environment but regretfully none were observed. Such signal absences are not uncommon in NQR spectroscopy and probably result from unfavorable nuclear relaxation conditions for those atoms.

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^{*}The eigenvalue equation for I = 9/2 given in ref. 13 was solved for values of the asymmetry parameter, η , in increments of 0.001 using a computer program written by L.F. Hamilton in this laboratory.